

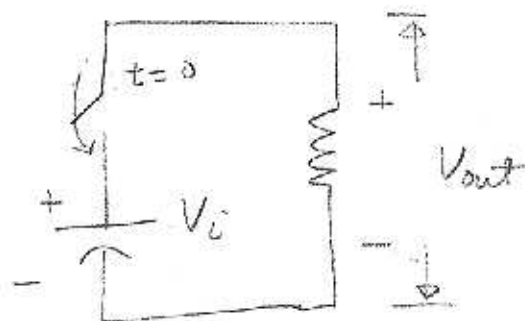
# Chapter 6 Continuous-Time Signal

## Analysis

6.1

A signal is called a deterministic signal if it can be described without any uncertainty.

Example:

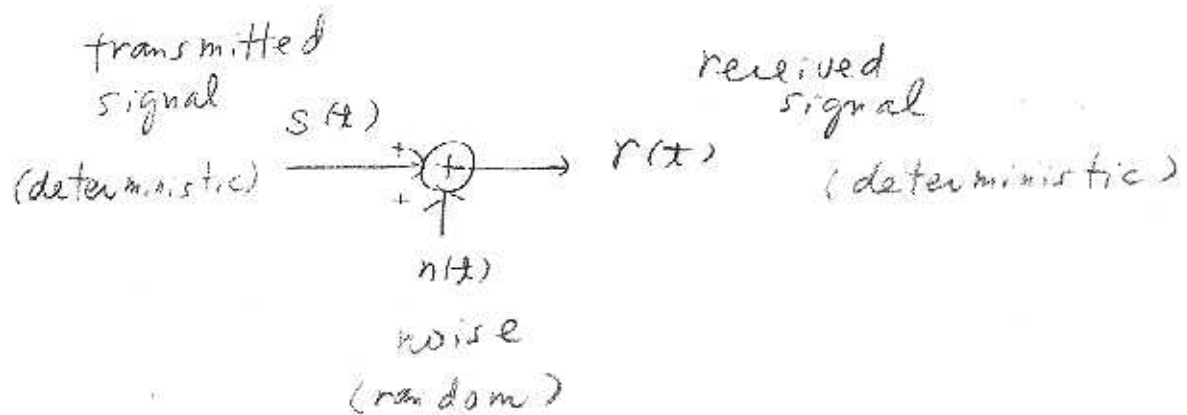


The voltage across the resistor is

$V_{out} = V_i g(x)$ .  $V_{out}$  is a deterministic signal.

On the contrary, a signal is called a random signal if it cannot be described with certainty before it actually occurs.

Example: The noise in the communication system.



### 6.1.1 Periodic and Aperiodic Signals

A function is said to be periodic with period  $P$  if  $f(t) = f(t+P)$  for all  $t$ . If  $f(t)$  is periodic,

then

$$f(t) = f(t+P) = f(t+2P) = \dots = f(t+nP)$$

for all  $t$  and for every integer  $n$ .

If a function is not periodic, it is non periodic or aperiodic function

If  $f_i(t)$  are periodic with period  $P_i$ ,  $i=1,2$ , then  $f_1(t) + f_2(t)$  is periodic with period  $P = n_1 P_1 = n_2 P_2$ ,  $n_1, n_2$  are two positive integers.  $f_1(t) + f_2(t)$  is aperiodic if no such  $n_1, n_2$  pair can be found.

Example. Which of the following are periodic functions? Find their fundamental frequencies and periods.

(a)  $\sin(3t) + \sin(\pi t)$

(b)  $\sin(0.8t) + \pi \cos(2t) - 2/\sin(1.2t)$

Answer:

(a)  $\sin(3(t+P_1)) = \sin(3t)$

$$3P_1 = 2\pi \Rightarrow P_1 = \frac{2\pi}{3}$$

$$\sin(\pi(t+P_2)) = \sin(\pi t)$$

$$\pi P_2 = 2\pi \Rightarrow P_2 = 2$$

$$n_1 P_1 = n_2 P_2$$

$$\Rightarrow \frac{n_1}{n_2} = \frac{P_2}{P_1} = \frac{3}{\pi}$$

No such integer  $n_1, n_2$  pair exists

$\therefore \sin(3t) + \sin(\pi t)$  is aperiodic

$$(b) \quad \sin(0.8(x + P_1)) = \sin(0.8x)$$

$$0.8 P_1 = 2\pi \Rightarrow P_1 = \frac{5}{2} \pi$$

$$\pi \cos(2(x + P_2)) = \pi \cos(2x)$$

$$2P_2 = 2\pi \Rightarrow P_2 = \pi$$

$$-21 \sin(1.2(x + P_3)) = -21 \sin(1.2x)$$

$$1.2 P_3 = 2\pi \Rightarrow P_3 = \frac{5}{3} \pi$$

$$P_1 = n_1 P_1 = n_2 P_2 = n_3 P_3$$

$$= \frac{5}{2} \pi n_1 = \pi n_2 = \frac{5}{3} \pi n_3$$

$$n_1 = 2$$

$$n_2 = 5$$

$$n_3 = 3$$

$$\left. \begin{array}{l} n_1 = 2 \\ n_2 = 5 \\ n_3 = 3 \end{array} \right\} \Rightarrow P = 5\pi \text{ radians}$$

$$\text{fundamental frequency} = \frac{2\pi}{P} = \frac{2\pi}{5\pi} = 0.4 \text{ Hz}$$

## 6.1.2 Energy and Power Signals

Energy provided by the signal  $f(x)$  over the time interval  $[t_0, t_1]$  is

$$\int_{t_0}^{t_1} |f(x)|^2 dx = \int_{t_0}^{t_1} f(x) f(x)^* dx$$

\* Energy signal: A signal  $f(x)$  is called an energy signal if its total energy in  $(-\infty, \infty)$  is finite, that is

$$E_{\infty} := \int_{-\infty}^{\infty} |f(x)|^2 dx < \infty$$

The average power of a signal over  $[t_1, t_2]$  is defined as

$$P_{av} := \frac{1}{t_2 - t_1} \int_{t_1}^{t_2} |f(x)|^2 dx$$

\* Power signal: A signal is called a power signal if the average power over  $(-\infty, \infty)$

$$P_{\infty} := \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T |f(x)|^2 dx$$

is nonzero and finite.

Example: The signal  $e^{-2t}$  is energy signal or power signal?

Answer:

$$\begin{aligned} E_{\infty} &= \int_{-\infty}^{\infty} |e^{-2t}|^2 dt \\ &= \int_{-\infty}^{\infty} e^{-4t} dt \\ &= -\frac{e^{-4t}}{4} \Big|_{-\infty}^{\infty} = 0 + \frac{e^{-4t}}{4} \Big|_{t \rightarrow -\infty} \\ &= \infty \end{aligned}$$

Hence it is not energy signal.

$$\begin{aligned} P_{\infty} &:= \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T e^{-4t} dt \\ &= \lim_{T \rightarrow \infty} \frac{1}{2T} \left[ -\frac{e^{-4t}}{4} + \frac{e^{4t}}{4} \right] \\ &= \lim_{T \rightarrow \infty} \left( -\frac{e^{-4T}}{8T} \right) + \lim_{T \rightarrow \infty} \left( \frac{e^{4T}}{8T} \right) \\ &= 0 + \lim_{T \rightarrow \infty} \frac{4e^{4T}}{8} = \infty \end{aligned}$$

it is not a power signal either.

### 6.1.3 Orthogonality of Complex Exponentials

$$e^{j\omega t} = \cos(\omega t) + j \sin(\omega t)$$

Consider the set of complex exponentials

$$\phi_m(t) := e^{jm\omega_0 t}, \text{ where } m = 0, \pm 1, \pm 2, \dots$$

For each integer  $m$ , the function  $e^{jm\omega_0 t}$  is a periodic function with fundamental frequency  $m\omega_0$  and fundamental period

$$P_m := \frac{2\pi}{m\omega_0}. \text{ If we define } P := \frac{2\pi}{\omega_0}$$

then  $P = mP_m$  and any time interval of length  $P$  contains  $m$  complete cycles of  $e^{jm\omega_0 t}$ .

Now we show

$$\begin{aligned} \int_{t_0}^{t_0+P} e^{jm\omega_0 t} dt &= \int_{t_0}^{t_0+P} [\cos(m\omega_0 t) + j \sin(m\omega_0 t)] dt \\ &= \begin{cases} P, & \text{if } m=0 \\ 0, & \text{if } m \neq 0 \end{cases} \end{aligned}$$

It implies

$$\int_{\langle P \rangle} \phi_m(t) \phi_n^*(t) dt := \int_{t_0}^{t_0+P} e^{j(m-n)\omega_0 t} dt$$
$$= \begin{cases} P, & \text{if } m=n \\ 0, & \text{if } m \neq n \end{cases}$$

This is called the orthogonality property of the set  $\phi_m(t)$

Consider a signal consisting of a linear combination of  $\phi_m(t)$  such that

$$f(t) = \sum_{m=-\infty}^{\infty} C_m e^{j m \omega_0 t} \quad \text{where } C_m \text{ are}$$

complex constants. It is clear that  $\omega_0$  is the GCD of all  $m \omega_0$ . Thus, for every set of  $C_m$ ,  $f(t)$  is periodic with fundamental period  $\omega_0$ . Conversely, every periodic function  $f(t)$  with fundamental period  $\omega_0$  can be represented by such a linear combination for some set of  $C_m$ .



## 6.2 Fourier Series of Periodic Functions

Consider a periodic function  $f(x)$  with fundamental period  $P$  and fundamental frequency  $\omega_0 = \frac{2\pi}{P}$ . We show that  $f(x)$  can be expressed as

$$f(x) = \sum_{m=-\infty}^{\infty} C_m e^{jm\omega_0 x} \quad \dots (X)$$

with

$$C_m = \frac{1}{P} \int_{t_0}^{t_0+P} f(x) e^{-jm\omega_0 t} dt$$

$$=: \frac{1}{P} \int_{\langle P \rangle} f(x) e^{-jm\omega_0 t} dt$$

for  $m = 0, \pm 1, \pm 2, \pm 3, \dots$  and for any arbitrary  $t_0$ .

Equation (X) is called the complex exponential Fourier Series.

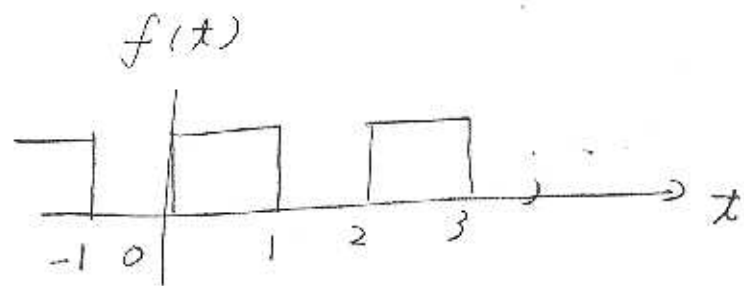
Proof:

$$\int_{t_0}^{t_0+P} f(x) e^{-jn\omega_0 t} dt = \sum_{m=-\infty}^{\infty} \int_{t_0}^{t_0+P} C_m e^{jm\omega_0 t} e^{-jn\omega_0 t} dt$$

$$= \sum_{m=-\infty}^{\infty} C_m \int_{t_0}^{t_0+P} e^{j(m-n)\omega_0 t} dt = PC_n$$

$$\therefore C_n = \frac{1}{p} \int_{t_0}^{t_0+p} f(t) e^{-jn\omega_0 t} dt, \quad n=0, \pm 1, \pm 2, \dots$$

Example :



$f(x)$  is depicted as above. Find the Fourier series.

Answer :

$$p = 2 \Rightarrow \omega_0 = \frac{2\pi}{2} = \pi$$

$$f(x) = \sum_{m=-\infty}^{\infty} C_m e^{jm\pi x}$$

$$\text{with } C_m = \frac{1}{2} \int_0^2 f(x) e^{-jm\pi x} dx$$

$$= \frac{1}{2} \int_0^1 e^{-jm\pi x} dx$$

$$= \frac{1}{2} \left. \frac{e^{-jm\pi x}}{-jm\pi} \right|_0^1 = \frac{e^{-jm\pi} - 1}{-2jm\pi}$$

$$f(x) = \sum_{m=-\infty}^{\infty} \left[ \frac{1 - e^{-jm\pi}}{2jm\pi} e^{jm\pi x} \right]$$

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Fourier Series

## 6.2.1 Discrete Frequency Spectrum

The discrete frequency spectrum is defined as the set of  $\{C_m\}$  and generally a complex-valued function of  $m$  or  $m\omega_0$ .

Let  $C_m = \alpha_m + j\beta_m$ . Then the magnitude and phase of  $C_m$  are

$$|C_m| := \sqrt{\alpha_m^2 + \beta_m^2}$$

$$\text{and } \angle C_m := \tan^{-1} \frac{\beta_m}{\alpha_m}$$

If  $f(x)$  is a real-valued function,

that is,  $f^*(x) = f(x)$ , then we have

$$C_m^* = \left[ \frac{1}{P} \int_{t_0}^{t_0+P} f(x) e^{-jm\omega_0 t} dt \right]^*$$

$$= \frac{1}{P} \int_{t_0}^{t_0+P} f(x) e^{jm\omega_0 t} dt$$

$$= C_{-m}$$

Thus the discrete frequency spectrum of real  $f(x)$  has the property

$$C_m^* = C_{-m}$$

This is called conjugate symmetry.

$$C_m^* = \alpha_m - j\beta_m = C_{-m} = \alpha_{-m} + j\beta_{-m}$$

$$\therefore \alpha_m = \alpha_{-m}, \quad -\beta_m = \beta_{-m}$$

Consequently,  $|C_m| = |C_{-m}|$  and

$$\angle C_m = -\angle C_{-m}$$

X: Even and Odd Real functions

If  $f(x)$  is a real and even function of  $t$  then  $C_m$  is a real and even function of  $m$ .

If  $f(x)$  is real and odd function of  $t$

then  $C_m$  is imaginary and odd function of  $m$ .

This can be established from

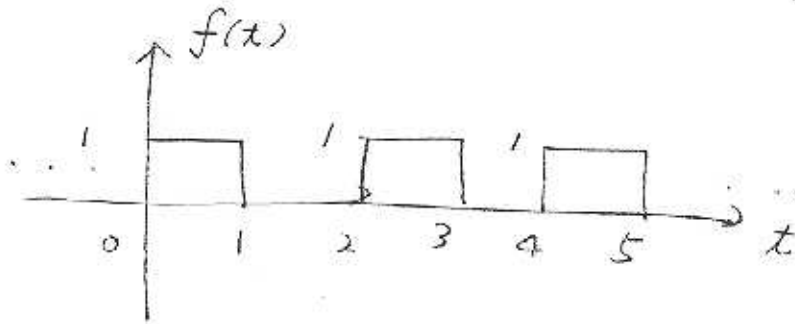
$$\begin{aligned} C_m &= \frac{1}{P} \int_{t_0}^{t_0+P} f(x) e^{-jm\omega_0 t} dt \\ &= \frac{1}{P} \int_{t_0}^{t_0+P} f(x) [\cos(m\omega_0 t) - j \sin(m\omega_0 t)] dt \end{aligned}$$

If  $f(x)$  is real, then

$$\operatorname{Re}\{C_m\} = \frac{1}{P} \int_{t_0}^{t_0+P} f(x) \cos(m\omega_0 t) dt$$

$$\text{and } \text{Im}\{C_m\} = -\frac{1}{P} \int_{t_0}^{t_0+P} f(t) \sin(m\omega_0 t) dt$$

Example: What is the discrete frequency spectrum for  $f(t)$  depicted as below?



$$C_m = \frac{1}{2} \int_0^1 e^{-jm\pi t} dt$$

$$= \frac{e^{-jm\pi} - 1}{-2jm\pi} = \frac{e^{-j\frac{m\pi}{2}} (e^{j\frac{m\pi}{2}} - e^{-j\frac{m\pi}{2}})}{2jm\pi}$$

$$= \frac{e^{-j\frac{m\pi}{2}} \sin\left(\frac{m\pi}{2}\right)}{m\pi}$$

$$|C_m| = \frac{|\sin\left(\frac{m\pi}{2}\right)|}{|m|\pi}$$

$$\angle C_m = -\frac{m\pi}{2} + \angle \sin\left(\frac{m\pi}{2}\right) - \angle m$$

\* Parseval's Formula:

$$\begin{aligned} P_{AV} &= \frac{1}{P} \int_{t_0}^{t_0+P} f(x) f^*(x) dt = \frac{1}{P} \int_{t_0}^{t_0+P} |f(x)|^2 dt \\ &= \sum_{m=-\infty}^{\infty} C_m C_m^* = \sum_{m=-\infty}^{\infty} |C_m|^2 \end{aligned}$$

Proof:

$$\begin{aligned} P_{AV} &:= \frac{1}{P} \int_{t_0}^{t_0+P} f(x) f^*(x) dt \\ &= \frac{1}{P} \int_{t_0}^{t_0+P} \left( \sum_{m=-\infty}^{\infty} C_m e^{jm\omega_0 t} \right) f^*(x) dt \\ &= \sum_{m=-\infty}^{\infty} C_m \left[ \frac{1}{P} \int_{t_0}^{t_0+P} f(x) e^{-jm\omega_0 t} dt \right]^* \\ &= \sum_{m=-\infty}^{\infty} C_m C_m^* \end{aligned}$$

#### 6.4. Fourier Transform

For any aperiodic function  $f(x)$ ,

the period  $P$  can be deemed as infinity.

Since the Fourier Series pair is given as

$$C_m = \frac{1}{P} \int_{-P/2}^{P/2} f(x) e^{-jm\omega_0 t} dt$$

$$\text{and } f(x) = \frac{1}{P} \sum_{m=-\infty}^{\infty} C_m P e^{jm\omega_0 t}$$

We define  $F(m\omega_0) := PC_m$

$$F(m\omega_0) := C_m P = \int_{-P/2}^{P/2} f(x) e^{-jm\omega_0 x} dx$$
$$= \int_{-P/2}^{P/2} f(x) e^{-j\omega x} dx$$

and  $f(x) = \frac{1}{P} \sum_{m=-\infty}^{\infty} PC_m e^{jm\omega_0 x}$

$$= \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} F(m\omega_0) e^{jm\omega_0 x}$$

Since  $f(x)$  is aperiodic,  $P \rightarrow \infty$ .

$$F(\omega) = \int_{-\infty}^{\infty} f(x) e^{-j\omega x} dx$$

$$f(x) = \lim_{\omega_0 \rightarrow 0} \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} F(m\omega_0) e^{jm\omega_0 x}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega x} d\omega$$

This is the Fourier Transform Pair.

$$F(\omega) = \int_{-\infty}^{\infty} f(x) e^{-j\omega x} dx$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega x} d\omega$$

Sufficient conditions for  $f(x)$  to have a Fourier transform are

1.  $f(x)$  is absolutely integrable over  $(-\infty, \infty)$ , that is  $\int_{-\infty}^{\infty} |f(x)| dt \leq M < \infty$  for some finite constant  $M$
2.  $f(x)$  has a finite number of discontinuities and a finite number of minima and maxima in every finite time interval.

Example: Determine the Fourier Transform for  $f(x) = e^{-a|x|}$ ,  $a > 0$  and  $a$  is real.

Answer:

$$\begin{aligned} F(\omega) &= \int_{-\infty}^{\infty} f(x) e^{-j\omega x} dt \\ &= \int_{-\infty}^0 e^{+ax} e^{-j\omega x} dt \\ &\quad + \int_0^{\infty} e^{-ax} e^{-j\omega x} dt \\ &= \frac{e^{(a-j\omega)x}}{a-j\omega} \Big|_{-\infty}^0 + \frac{e^{-(a+j\omega)x}}{-(a+j\omega)} \Big|_0^{\infty} \end{aligned}$$



$$= \frac{1}{a-j\omega} - \frac{1}{-a-j\omega} = \frac{a+j\omega + a-j\omega}{a^2 + \omega^2} = \frac{2a}{a^2 + \omega^2}$$

$F(\omega)$  is usually called the frequency spectrum.

$$A(\omega) := |F(\omega)| = \sqrt{\operatorname{Re}\{F(\omega)\}^2 + \operatorname{Im}\{F(\omega)\}^2}$$

is defined as the amplitude spectrum

and

$$\theta(\omega) := \angle F(\omega) = \tan^{-1} \frac{\operatorname{Im}\{F(\omega)\}}{\operatorname{Re}\{F(\omega)\}}$$

is defined as the phase spectrum.

If  $f(t)$  is real, then  $A(\omega) = A(-\omega)$  and

$$\theta(\omega) = -\theta(-\omega)$$

If  $f(t)$  is real and even function of  $t$

( $f(t) = f(-t)$ ), then  $F(\omega)$  is a real and

even function of  $\omega$ . If  $f(t)$  is a real and

odd function of  $t$  ( $f(t) = -f(-t)$ ), then

$F(\omega)$  is an imaginary and odd function of  $\omega$ .

\* The Fourier Transform is a linear operator, that is,

$$\begin{aligned} \mathcal{F} [ \alpha_1 f_1(x) + \alpha_2 f_2(x) ] \\ &= \alpha_1 \mathcal{F} [ f_1(x) ] + \alpha_2 \mathcal{F} [ f_2(x) ] \\ &= \alpha_1 F_1(\omega) + \alpha_2 F_2(\omega) \end{aligned}$$

where  $\alpha_1, \alpha_2$  are two arbitrary complex-valued constants, and

$$\begin{aligned} \mathcal{F} [ f_i(x) ] &= F_i(\omega) \\ &= \int_{-\infty}^{\infty} f_i(t) e^{-j\omega t} dt \end{aligned}$$

6.4.1 From the Laplace Transform to the Fourier Transform

The Fourier Transform can be established from the Laplace Transform evaluated at the pure imaginary axis, namely  $s = \pm j\omega$

Consider an absolutely integrable function  $f(t)$  defined over  $(-\infty, \infty)$ . (ITF Fourier Transform exists.)

$$f(t) = f_-(t) + f_+(t) = \overbrace{f(t)g(-t)}^{f_-(t)} + \overbrace{f(t)g(t)}^{f_+(t)} + f(t)g(t),$$

where  $g(t)$  is redefined

$$g(t) := \begin{cases} 1, & t > 0 \\ \frac{1}{2}, & t = 0 \\ 0, & t < 0 \end{cases}$$

$$\mathcal{F}[f(t)] = \mathcal{F}[f_+(t)] + \mathcal{F}[f_-(t)]$$

$$F_+(\omega) = \mathcal{F}[f_+(t)]$$

$$= \int_{-\infty}^{\infty} f_+(t) e^{-j\omega t} dt$$

$$= \int_{0^-}^{\infty} f_+(t) e^{-j\omega t} dt$$

$$\text{Let } \bar{F}_+(s) = \mathcal{L}[f_+(t)]$$

$$= \int_{0^-}^{\infty} f_+(t) e^{-st} dt$$

$$\therefore F_+(\omega) = \bar{F}_+(s) \big|_{s=j\omega}$$

$$F_+(w) = \mathcal{F} [f(x)g(x)] = \mathcal{L} [f(x)g(x)] \Big|_{s=jw}$$

Similarly

$$F_-(w) = \int_{-\infty}^{\infty} f_-(x) e^{-jwx} dx$$

$$= \int_{-\infty}^{0^+} f_-(x) e^{-jwx} dx$$

$$\text{let } \tau = -x$$

$$F_-(w) = \int_{0^-}^{\infty} f_-(\tau) e^{jw\tau} d\tau$$

$$\text{Let } \bar{F}_-(s) = \mathcal{L} [f_-(\tau)]$$

$$= \int_{0^-}^{\infty} f_-(\tau) e^{-s\tau} d\tau$$

$$\mathcal{F} [f_-(x)] = F_-(w) = \mathcal{L} [f_-(\tau)] \Big|_{s=-jw}$$

$$F_-(w) = \mathcal{F} [f(x)g(-x)] = \mathcal{L} [f(-x)g(x)] \Big|_{s=jw}$$

$$F(w) = F_+(w) + F_-(w)$$

$$= \bar{F}_+(s) \Big|_{s=jw} + \bar{F}_-(s) \Big|_{s=-jw}$$

Example : Find the Fourier Transform  
for  $f(t) = \begin{cases} e^{-2t} & , t > 0 \\ -e^{2t} & , t < 0 \end{cases}$

Answer :

$$f_+(t) = e^{-2t}, \quad t > 0$$

$$f_-(t) = -e^{2t}, \quad t < 0$$

$$\bar{F}_+(s) = \mathcal{L}[f_+(t)] = \frac{1}{s+2}$$

$$\bar{F}_-(s) = \mathcal{L}[f_-(t)] = -\frac{1}{s+2}$$

$$F_+(\omega) = \left. \frac{1}{s+2} \right|_{s=j\omega} = \frac{1}{j\omega+2}$$

$$F_-(\omega) = \left. -\frac{1}{s+2} \right|_{s=-j\omega} = -\frac{1}{-j\omega+2} = \frac{1}{j\omega-2}$$

$$F(\omega) = F_+(\omega) + F_-(\omega)$$

$$= \frac{1}{j\omega+2} + \frac{1}{j\omega-2} = \frac{-2j\omega}{4+\omega^2}$$

Since  $f(t)$  is odd function,

$F(\omega)$  is a pure imaginary and  
odd function.

## 6.4.2 Fourier Transform of Periodic Functions

The impulse  $F(\omega) = \delta(\omega - \omega_0)$  will cause a complex exponential function

$$f(t) \left[ \begin{array}{c} \xrightarrow{\mathcal{F}} \\ \xleftarrow{\mathcal{F}^{-1}} \end{array} F(\omega) \right]$$

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega - \omega_0) e^{j\omega t} d\omega$$

$$= \frac{1}{2\pi} e^{j\omega_0 t}$$

$\therefore$  We can conclude

$$e^{j\omega_0 t} \xrightarrow{\mathcal{F}} 2\pi \delta(\omega - \omega_0)$$
$$\xleftarrow{\mathcal{F}^{-1}}$$

Example: Find the Fourier Transforms for  $\sin(\omega_0 t)$  and  $\cos(\omega_0 t)$

$$\text{Answer: } \mathcal{F}[\sin(\omega_0 t)] = \mathcal{F}\left[\frac{e^{j\omega_0 t} - e^{-j\omega_0 t}}{2j}\right]$$

$$= \frac{1}{2j} \mathcal{F}[e^{j\omega_0 t}] - \frac{1}{2j} \mathcal{F}[e^{-j\omega_0 t}]$$

$$= \frac{\pi}{j} [\delta(\omega - \omega_0) - \delta(\omega + \omega_0)]$$

$$\begin{aligned} \mathcal{F}[\cos(\omega_0 t)] &= \mathcal{F}\left[\frac{1}{2}(e^{j\omega_0 t} + e^{-j\omega_0 t})\right] \\ &= \frac{1}{2}\mathcal{F}[e^{j\omega_0 t}] + \frac{1}{2}\mathcal{F}[e^{-j\omega_0 t}] \\ &= \pi[\delta(\omega - \omega_0) + \delta(\omega + \omega_0)] \end{aligned}$$

If  $f(x)$  is periodic function, then

$$f(x) = \sum_{m=-\infty}^{\infty} C_m e^{jm\omega_0 t} \quad (\text{Fourier Series})$$

$$F(\omega) = \mathcal{F}[f(x)] = \sum_{m=-\infty}^{\infty} 2\pi C_m \delta(\omega - m\omega_0)$$

Example: Find the Fourier Transform for

a function  $f(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT)$

Answer:  $f(t) = \sum_{m=-\infty}^{\infty} C_m e^{jm\omega_0 t}$

where  $\omega_0 = \frac{2\pi}{T}$

$$C_m = \frac{1}{T} \int_{-T/2}^{T/2} \delta(t) e^{-jm\omega_0 t} dt$$

$$= \frac{1}{T}$$

$$\therefore f(t) = \sum_{m=-\infty}^{\infty} \frac{1}{T} e^{jm\omega_0 t} \quad (\text{Fourier Series})$$

$$\text{Then } F(\omega) = \sum_{m=-\infty}^{\infty} \frac{2\pi}{T} \delta(\omega - m\omega_0)$$

Table 6.1 lists some Fourier Transform pairs!

#### 6.4.4 Spectrum density and Energy

Consider an absolutely integrable function  $f(x)$ , let  $F(\omega)$  be its Fourier transform or frequency spectrum. The total energy

of  $f(x)$  is

$$\begin{aligned} \text{total energy} &:= E = \int_{-\infty}^{\infty} |f(x)|^2 dx \\ &= \int_{-\infty}^{\infty} f(x) f^*(x) dx \end{aligned}$$

$$E = \int_{-\infty}^{\infty} |f(x)|^2 dx = \int_{-\infty}^{\infty} f(x) \left[ \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega x} d\omega \right] dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F^*(\omega) \left[ \int_{-\infty}^{\infty} f(x) e^{-j\omega x} dx \right] d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F^*(\omega) F(\omega) d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(\omega)|^2 d\omega$$



## 6.5 Convolution

The zero-state response of every LTI continuous time system can be described as

$$y(t) = \int_{-\infty}^{\infty} h(t-\tau) u(\tau) d\tau$$

$$\mathcal{F}[y(t)] = Y(\omega) = \int_{-\infty}^{\infty} y(t) e^{-j\omega t} dt$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(t-\tau) u(\tau) d\tau e^{-j\omega(t-\tau)} e^{-j\omega\tau} dt$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(t-\tau) e^{-j\omega(t-\tau)} dt u(\tau) e^{-j\omega\tau} d\tau$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(t') e^{-j\omega t'} dt' u(\tau) e^{-j\omega\tau} d\tau$$

$$= H(\omega) \underbrace{\int_{-\infty}^{\infty} u(\tau) e^{-j\omega\tau} d\tau}_{U(\omega)}$$

$$= H(\omega) U(\omega)$$

Example: Find the Fourier Transform for  $y(t)$

$$= h(t) \otimes u(t), \text{ where } u(t) = g(t)$$

$$\text{and } h(t) = e^{-t} g(t)$$

Answer:

$$g(t) = \frac{1}{2} + \frac{1}{2} \operatorname{sgn}(t)$$

$$\text{where } \operatorname{sgn}(t) = \begin{cases} 1 & , t > 0 \\ \frac{1}{2} & , t = 0 \\ -1 & , t < 0 \end{cases}$$

$$\mathcal{F}[\operatorname{sgn}(t)] = \lim_{a \rightarrow 0} \left[ - \int_{-\infty}^0 e^{at} e^{-j\omega t} dt \right.$$

$$\left. + \int_0^{\infty} e^{-at} e^{-j\omega t} dt \right]$$

$$= \lim_{a \rightarrow 0} \left[ \frac{-1}{a-j\omega} + \frac{1}{a+j\omega} \right]$$

$$= \frac{2}{j\omega}$$

$$\mathcal{F}\left[\frac{1}{2}\right] = \mathcal{F}\left[\frac{1}{2} e^{j0t}\right] = \pi \delta(\omega - 0) \\ = \pi \delta(\omega)$$

$$\therefore \mathcal{F}[g(t)] = Q(\omega) = \frac{1}{j\omega} + \pi \delta(\omega)$$

$$H(\omega) = \mathcal{L}[e^{-t} g(t)] \Big|_{s=j\omega} + 0$$

$$= \frac{1}{s+1} \Big|_{s=j\omega} = \frac{1}{j\omega+1}$$

$$Y(\omega) = \left[ \frac{1}{j\omega} + \pi \delta(\omega) \right] \left( \frac{1}{j\omega+1} \right)$$

## 6.6 Properties of Fourier Transform

(i) Modulation in time is equivalent to shifting in Frequency:

$$\mathcal{F} [ f(x) e^{j\omega_0 x} ] = F(\omega - \omega_0)$$

Proof:

$$\begin{aligned} \mathcal{F} [ f(x) e^{j\omega_0 x} ] &= \int_{-\infty}^{\infty} f(x) e^{j\omega_0 x} e^{-j\omega x} dx \\ &= \int_{-\infty}^{\infty} f(x) e^{-j(\omega - \omega_0)x} dx = F(\omega - \omega_0) \quad \# \end{aligned}$$

(ii) Time-shifting Property:

$$\mathcal{F} [ f(x - t_0) ] = e^{-j\omega t_0} F(\omega)$$

Proof:

$$\begin{aligned} \mathcal{F} [ f(x - t_0) ] &= \int_{-\infty}^{\infty} f(x - t_0) e^{-j\omega x} dx \\ &= \int_{-\infty}^{\infty} f(x') e^{-j\omega(x'+t_0)} dx' \quad x' = x - t_0 \\ &= e^{-j\omega t_0} \underbrace{\int_{-\infty}^{\infty} f(x') e^{-j\omega x'} dx'}_{F(\omega)} \\ &= e^{-j\omega t_0} F(\omega) \quad \# \end{aligned}$$

(iii) Duality Property:

$$\text{If } F(\omega) = \mathcal{F} [ f(x) ]$$

$$\text{then } f(-\omega) = \frac{1}{2\pi} \mathcal{F} [ F(x) ]$$

Proof:

$$F(\omega) = \int_{-\infty}^{\infty} f(x) e^{-j\omega x} dx$$

change  $x$  to  $x'$ , change  $\omega$  to  $\omega'$

$$F(x') = \int_{-\infty}^{\infty} f(x) e^{-j\omega' x'} dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} 2\pi f(x) e^{-j\omega' x'} dx$$

$$\text{let } \omega' = -\omega$$

$$F(x') = \frac{1}{2\pi} \int_{-\infty}^{\infty} 2\pi f(-\omega') e^{j\omega' x'} d\omega'$$

$$\Rightarrow F(x') = \mathcal{F}^{-1} [2\pi f(-\omega')]$$

$$\text{or } \mathcal{F} [F(x')] = 2\pi f(-\omega') \quad \#$$

(iv) Time-Scaling Property:

$$\text{If } F(\omega) = \mathcal{F} [f(x)]$$

$$\text{then } \mathcal{F} [f(\alpha x)] = \frac{1}{|\alpha|} F\left(\frac{\omega}{\alpha}\right)$$

$$\text{Proof: } \mathcal{F} [f(\alpha x)] = \int_{-\infty}^{\infty} f(\alpha x) e^{-j\omega x} dx$$

$$= \int_{-\infty}^{\infty} f(x') e^{-j\frac{\omega}{\alpha} x'} dx', \quad \alpha > 0$$
$$= \begin{cases} \frac{1}{\alpha} \int_{-\infty}^{\infty} f(x') e^{-j\frac{\omega}{\alpha} x'} dx', & \alpha > 0 \\ -\frac{1}{\alpha} \int_{-\infty}^{\infty} f(x') e^{-j\frac{\omega}{\alpha} x'} dx', & \alpha < 0 \end{cases}$$

$$= \frac{1}{|\alpha|} F\left(\frac{\omega}{\alpha}\right) \quad \#$$

(V) Complex convolution:

$$\mathcal{F}[f_1(x) f_2(x)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} F_1(\beta) F_2(\omega - \beta) d\beta$$

where  $F_1(\omega) = \mathcal{F}[f_1(x)]$  and  $F_2(\omega) = \mathcal{F}[f_2(x)]$

Proof:  $\mathcal{F}[f_1(x) f_2(x)]$

$$= \int_{-\infty}^{\infty} f_1(x) f_2(x) e^{-j\omega x} dx$$

$$= \int_{-\infty}^{\infty} f_1(x) \left[ \frac{1}{2\pi} \int_{-\infty}^{\infty} F_2(\beta) e^{j\beta x} d\beta \right] e^{-j\omega x} dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \underbrace{\int_{-\infty}^{\infty} f_1(x) e^{-j(\omega - \beta)x} dx}_{F_1(\omega - \beta)} F_2(\beta) d\beta$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F_1(\omega - \beta) F_2(\beta) d\beta$$

$$\text{Since } \mathcal{F}[f_1(x) f_2(x)] = \mathcal{F}[f_2(x) f_1(x)]$$

$$\mathcal{F}[f_1(x) f_2(x)] = \mathcal{F}[f_2(x) f_1(x)]$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F_1(\omega - \beta) F_2(\beta) d\beta$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} F_1(\beta) F_2(\omega - \beta) d\beta \quad \#$$

Example: What is the energy

$$E = \int_{-\infty}^{\infty} |f(x)|^2 dx \quad \text{for } f(x) = \frac{-2jx}{4+x^2}$$

Answer: Since  $\int_{-\infty}^{\infty} |f(x)|^2 dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(\omega)|^2 d\omega$

$$\mathcal{F}[f(x)] = F(\omega)$$

From previous example,

$$f(x) = \begin{cases} e^{-2x}, & x > 0 \\ -e^{2x}, & x < 0 \end{cases} \quad \begin{matrix} \xrightarrow{\mathcal{F}} \\ \xleftarrow{\mathcal{F}^{-1}} \end{matrix}$$

$$F(\omega) = \frac{-2j\omega}{4+\omega^2}$$

According to Property (iii),

$$\mathcal{F}[F(x)] = \mathcal{F}\left[\frac{-2jx}{4+x^2}\right] = 2\pi f(-\omega)$$

$$= \begin{cases} 2\pi e^{-2\omega}, & \omega > 0 \\ -2\pi e^{-2\omega}, & \omega < 0 \end{cases}$$

$$E = \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(\omega)|^2 d\omega$$

$$= 2\pi \int_0^{\infty} e^{-4\omega} d\omega + 2\pi \int_{-\infty}^0 e^{-4\omega} d\omega$$

$\rightarrow \infty$